

# CRYSTAL BINDING Lecture 7

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# Crystal Binding (cont)

3.5.1 Energy of van der Waals Solid

**Convert from pair-wise interactions by summing over all pairs.** 

- pick an atom, label it 0,
- $\bullet$  let energy of interaction of this atom with neighbour i at a distance  $r_i$  be  $U(r_i)$
- total potential energy of the atom is

$$U_0 = \sum_i U(r_i)$$

- if there are N atoms altogether, each will have this same energy but interaction 0 i is the same as i 0
- so total energy is

$$U = \frac{N}{2} \sum_{i} U(r_i).$$

Separate structure (spatial arrangement) from scale (interatomic separation).

• write  $r_i = \rho_i r_0$ , where  $r_0$  is nearest neighbour distance,  $\rho_i$  is dimensionless

• then

$$U(r_i) = 4\epsilon \left[ \frac{1}{\rho_i^{12}} \left( \frac{\sigma}{r_0} \right)^{12} - \frac{1}{\rho_i^6} \left( \frac{\sigma}{r_0} \right)^6 \right]$$

• summing,

$$U = 2N\epsilon \left[ A_{12} \left( \frac{\sigma}{r_0} \right)^{12} - A_6 \left( \frac{\sigma}{r_0} \right)^6 \right],$$

where

$$A_n = \sum_i \frac{1}{\rho_i^n}.$$

• these lattice sums can be done for any structure.

Structure	$A_{12}$	$A_6$
FCC	12.12188	14.45392
HCP	12.13229	14.45489
BCC	9.11418	12.25330

Note:

- expect sum of  $1/r^n$  to converge rapidly for large n
- $A_{12}$  is dominated by the nearest neighbours (10 in FCC, HCP, 8 in BCC), but more distant neighbours affect  $A_6$

#### **3.5.2 Equilibrium Separation**

The equilibrium structure minimises the total energy:  $\partial U/\partial r_0 = 0$ .

$$\frac{\partial U}{\partial r_0} = -2N\epsilon \left[ 12A_{12} \frac{\sigma^{12}}{r_0^{13}} - 6A_6 \frac{\sigma^6}{r_0^7} \right],$$

which is zero when

$$\frac{r_0}{\sigma} = \left(\frac{2A_{12}}{A_6}\right)^{1/6}$$

 $\frac{r_0}{\sigma} = 1.09$  for FCC.

$$U = -\frac{A_6^2}{2A_{12}}\epsilon \text{ per atom.}$$

Typically about 0.01 to 0.1 eV per atom.

#### 3.5.3 Choice of Structure

Expect structure to form crystals which have lowest energy, i.e. largest cohesive energy. Strictly, Gibbs free energy,

$$G = U - TS + pV,$$

but assume T = 0 and p = 0. Neglect kinetic energy of atomic motion  $\frac{\text{SC BCC HCP FCC}}{A_6 8.4 12.25 14.45 14.45}$ 

in U.  $A_{12}$  6.2 9.11 12.13 12.12

$$\frac{U}{N\epsilon}$$
 -5.69 -8.24 -8.61 -8.62

Note how close FCC and HCP are in energy – but FCC is favoured.

#### **3.5.4 Bulk Modulus**

We know energy as a function of separation: need to express as function of volume. For FCC structure, cubic lattice parameter a, nearest-neighbour separation  $r_0 = a/\sqrt{2}$ . Cubic unit cell, volume  $a^3$ , contains 4 atoms, so

Volume per atom = 
$$\frac{a^3}{4}$$
  
=  $\frac{r_0^3}{\sqrt{2}}$ 

Now we could use

$$\frac{\partial}{\partial V} = \frac{\partial r_0}{\partial V} \frac{\partial}{\partial r_0}$$

but it's easier to substitute

$$r_0 = 2^{1/6} V^{1/3} N^{-1/3}$$

in

$$U = 2N\epsilon \left[ A_{12} \left( \frac{\sigma}{r_0} \right)^{12} - A_6 \left( \frac{\sigma}{r_0} \right)^6 \right],$$

to get

$$U = 2N\epsilon \left[ A_{12} \frac{\sigma^{12} N^4}{4V^4} - A_6 \frac{\sigma^6 N^2}{2V^2} \right]$$

and hence

$$\frac{\partial^2 U}{\partial V^2} = 2N\epsilon \left[ A_{12} \frac{20\sigma^{12}N^4}{4V^6} - A_6 \frac{6\sigma^6 N^2}{2V^4} \right]$$

**SO** 

$$B = N\epsilon \left[ A_{12} \frac{10\sigma^{12}N^4}{V^5} - A_6 \frac{6\sigma^6 N^2}{V^3} \right].$$

**But in equilibrium** 

$$r_0 = \sigma \left(\frac{2A_{12}}{A_6}\right)^{1/6},$$

**SO** 

$$V = N\sigma^3 \sqrt{\left(\frac{A_{12}}{A_6}\right)}$$

and

$$B = N\epsilon \left[ 10A_{12}\sigma^{12}N^4 \times \frac{1}{N^5\sigma^{15}} \left(\frac{A_6}{A_{12}}\right)^{5/2} - 6A_6\sigma^6N^2 \times \frac{1}{N^3\sigma^9} \left(\frac{A_6}{A_{12}}\right)^{3/2} \right],$$

which simplifies to

$$B = \frac{4A_6^{5/2}\epsilon}{A_{12}^{3/2}\sigma^3}.$$

### **3.6 Ionic Crystals**

The picture of an assembly of spherical ions is a good one:



(Theoretical calculations by Harker, checked against experiment)

#### 3.6.1 Ionic Radii and Packing

In general, cation  $M^+$  and anion  $X^-$  have different radii. We expect lowest energy if we have as many cations as possible around each anion, and we avoid anions touching anions. We know that for equal-sized spheres FCC gives high packing.



If we shrink the smaller ions, but keep the geometrical arrangement, eventually the larger ions will touch.



NaCl: X atoms touch if

$$\begin{aligned} r_{MX} &= r_M + r_X \\ r_{XX} &= \sqrt{2}r_{MX} \\ r_{XX} &\leq 2r_X \rightarrow \sqrt{2}(r_M + r_X) \leq 2r_X \\ \frac{r_X}{r_M} &\geq \frac{1}{\sqrt{2} - 1}. \end{aligned}$$



#### **CsCl:** X atoms touch if

$$r_{MX} = r_M + r_X$$

$$r_{XX} = \frac{2}{\sqrt{3}} r_{MX}$$

$$r_{XX} \le 2r_X \rightarrow \frac{2}{\sqrt{3}} (r_M + r_X) \le 2r_X$$

$$\frac{r_X}{r_M} \ge \frac{1}{\sqrt{3} - 1}.$$

Given a table of ionic radii, we can guess structures of compounds.

#### **3.6.2** Ionic Lattice Sums

For a pair of ions,

$$U_{ij} = \frac{q_i q_j e^2}{4\pi\epsilon_0 r_{ij}} + U_{\rm rep}(r_{ij}),$$

and summing as before gives

$$U = \frac{N}{2} \left[ -\alpha_M \frac{e^2}{4\pi\epsilon_0 r_0} + U_{\rm rep} \right]$$

 $\alpha_M$  is the *Madelung constant*, obtained by a lattice sum:

$$-\alpha_M = \sum_i \frac{q_0 q_i}{\rho_i}$$

#### 3.6.3 Linear Chain



#### Note:

- very slowly convergent
- only converges because it is an alternating series try to sum only effect of, say, positive ions and get infinity
- **Result:**  $\alpha_M = 2 \ln(2)$ .

#### **3.6.4** Three dimensions

Special mathematical tricks used to calculate Madelung constant.

- Evjen method: sum neutral regions, using increasingly large cubes and only counting half of charges on face centres, quarter of cube edges, eighth of cube corners
- Ewald method: trick involving real space and reciprocal space

Structure	coordination	$lpha_M$
	number	
CsCl	8	1.7627
NaCl	6	1.7476
Zinc blende (like GaAs)	4	1.6381
Wurtzite (hexagonal ZnS)	4	1.641
Higher coordination gives larger Ma	delung constan	nt.

#### **3.6.5** Ionic Structures

#### Structure will be that which minimises energy.



Energy increasingly negative as ions get closer – until like ions touch. Radius ratios (smaller  $r_s$  over larger  $r_l$ ) give good guidance. Similar radii favour close packed structures – very different radii give more open, lower-coordinated (and more covalent) structures.

## Summary

- Binding Energy, Equilibrium Separation
- Lattice Sums
- Coulomb sums
- Minimum-energy structures
- Elasticity

#### Next:

• Lattice vibrations